



Polarisation as a tool to study the  
Solar System and beyond

Action MP1104



# X-ray polarization in the lamp-post model of non-smooth black-hole accretion discs

**Michal Dovčiak**

Astronomical Institute  
Academy of Sciences of the Czech Republic, Prague

**R. W. Goosmann**

Strasbourg Astronomical  
Observatory

**F. Marin & V. Karas**

Astronomical Institute  
Prague

**G. Matt**

University  
Roma Tre

**F. Muleri**

IASF  
Rome

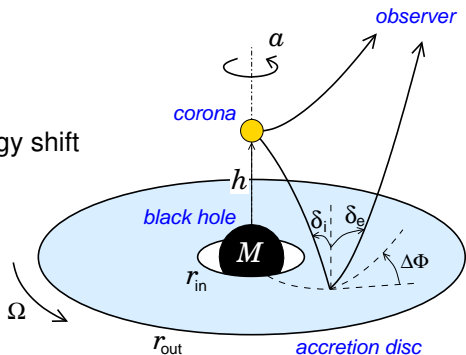
*X-ray polarisation in astrophysics: a window about to open?*

AlbaNova University Center, Stockholm, Sweden

25<sup>th</sup> – 28<sup>th</sup> August 2014

# Scheme of the lamp-post geometry

- ▶ **central black hole** → mass, spin
- ▶ **accretion disc**
  - Keplerian, geometrically thin, optically thick and neutral
- ▶ **compact corona**
  - isotropic and unpolarised power-law emission
  - static (or slow motion)
  - height, photon index
- ▶ **relativistic effects:**
  - Doppler and gravitational energy shift
  - light bending (lensing)
  - aberration (beaming)
- ▶ **references:**
  - Matt (1993)
  - Dovčiak, Muleri, Goosmann, Karas & Matt (2011)



# Motivation

- ▶ observational evidence of a rather compact X-ray source: variability, micro-lensing  $\rightarrow$  corona size of ten(s) of  $R_g = GM/c^2$
- ▶ many effects should be qualitatively similar with a simple lamp-post geometry
- ▶ base of an aborted jet?
- ▶ light bending scenario to explain variability in continuum versus line flux (e.g. Miniutti & Fabian, 2004)
- ▶ polarization in reflection in non-axisymmetric geometry should be significant
- ▶ reflection versus absorption scenarios for origin of the AGN X-ray spectral shape in 2–10 keV range (Marin et al, 2012, 2013)
- ▶ new polarimetric detectors for next generation X-ray missions have been developed and proposed (XEUS, IXO, NHXM, GEMS, XIPE)

# Stokes parameters at infinity

$$\Delta I(E) = \int_{\Sigma} dS \, G \, I_{\text{loc}}(E/g)$$

$$G = g^3 \ell \mu_e$$

$$\Delta Q(E) = \int_{\Sigma} dS \, G \, P_{\text{loc}}(E/g) \, I_{\text{loc}}(E/g) \cos 2[\chi_{\text{loc}}(E/g) + \psi]$$

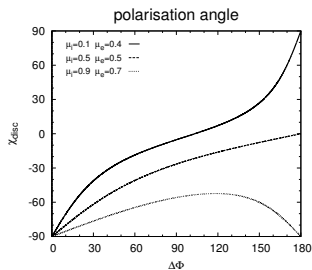
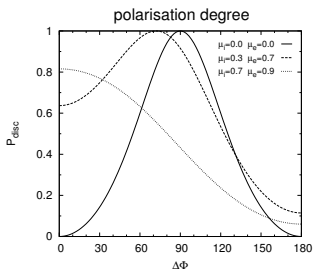
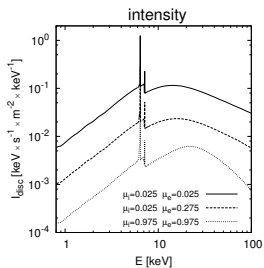
$$\Delta U(E) = \int_{\Sigma} dS \, G \, P_{\text{loc}}(E/g) \, I_{\text{loc}}(E/g) \sin 2[\chi_{\text{loc}}(E/g) + \psi]$$

$$P = \frac{\sqrt{(\Delta Q)^2 + (\Delta U)^2}}{\Delta I} \quad \tan 2\chi = \frac{\Delta U}{\Delta Q}$$

→  $I_{\text{loc}}$ ,  $P_{\text{loc}}$  and  $\chi_{\text{loc}}$  depend on local geometry of scattering

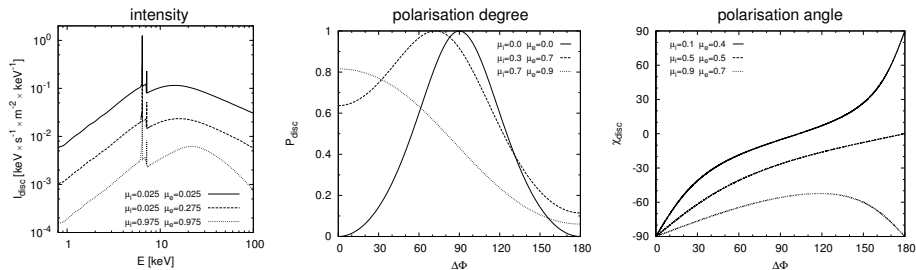


# Local emission



- ▶ flux → multiple Compton scattering and  $K\alpha$ ,  $K\beta$  fluorescence – NOAR
- ▶ polarization → single scattering approximation – Chandrasekhar (1960)
- ▶ irregular surface
  - ▶ to relax the assumption on such a well defined scattering geometry
  - ▶ on much smaller scale than the changes in relativistic effects
  - ▶ on much larger scale than the scale characterizing the radiative transfer
  - ▶ averaging of the local polarisation properties

# Local emission



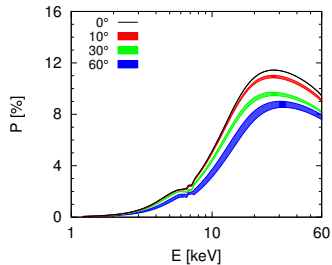
- ▶ **flux** → multiple Compton scattering and  $K\alpha$ ,  $K\beta$  fluorescence – NOAR
- ▶ **polarization** → single scattering approximation – Chandrasekhar (1960)
- ▶ **irregular surface**
  - ▶ defined by the angle of the surface with respect to the equatorial plane,  $\delta < \delta_{\text{max}} = 10^\circ, 30^\circ$  and  $60^\circ$
  - ▶ locally symmetric in azimuthal angle
  - ▶ shading is not taken into account
  - ▶ two probability distributions are explored,  $\mathcal{P}(\delta)$ :
    - linearly decreases to zero for  $\delta_{\text{max}}$
    - linearly increases to be doubled for  $\delta_{\text{max}}$

# Smooth disc

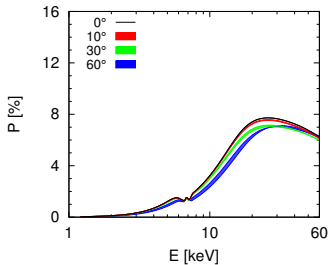
- importance of the local polarization properties
  - geometry of scattering (incident, emission and relative azimuthal angles)
  - source height, observer inclination and black hole spin
  - formation of additional depolarizing critical points
  - illumination pattern depends on height of the source
-

# Irregular disc surface – polarisation degree

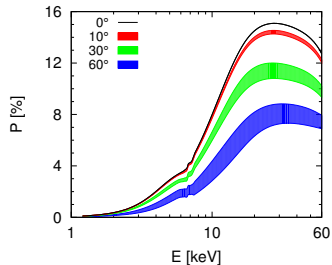
$a=0, \theta_0=30^\circ, h=3$



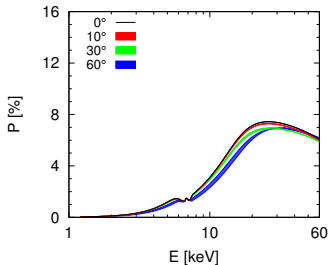
$a=0, \theta_0=30^\circ, h=15$



$a=1, \theta_0=30^\circ, h=3$

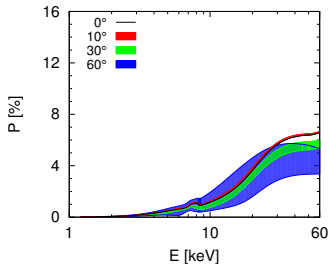


$a=1, \theta_0=30^\circ, h=15$

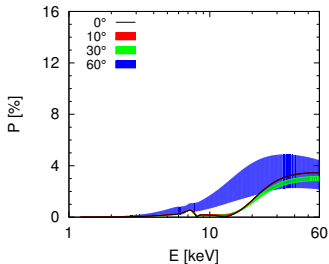


# Irregular disc surface – polarisation degree

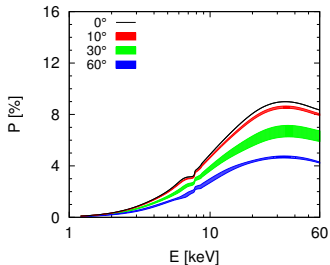
$a=0, \theta_0=60^\circ, h=3$



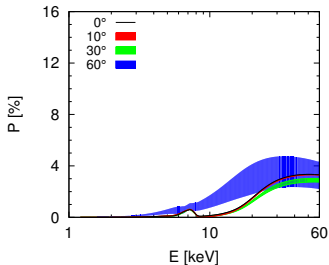
$a=0, \theta_0=60^\circ, h=15$



$a=1, \theta_0=60^\circ, h=3$

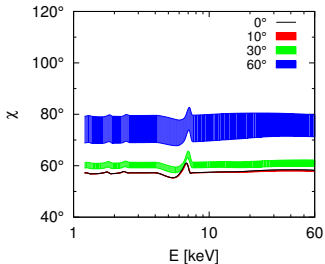


$a=1, \theta_0=60^\circ, h=15$

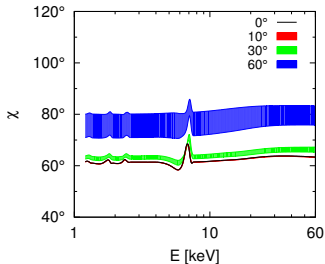


# Irregular disc surface – polarisation angle

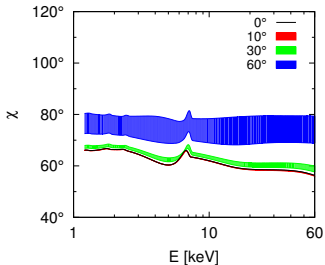
$a=0, \theta_0=30^\circ, h=3$



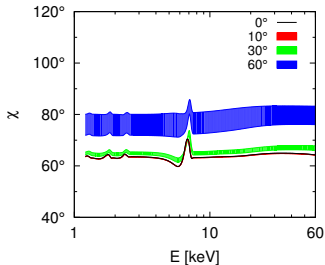
$a=0, \theta_0=30^\circ, h=15$



$a=1, \theta_0=30^\circ, h=3$

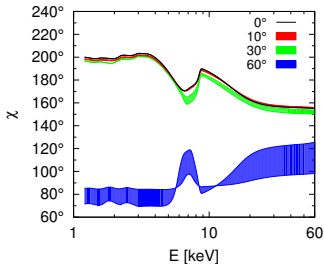


$a=1, \theta_0=30^\circ, h=15$

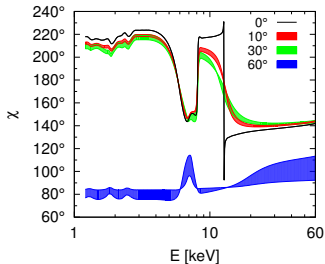


# Irregular disc surface – polarisation angle

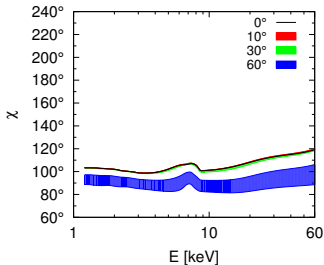
$a=0, \theta_0=60^\circ, h=3$



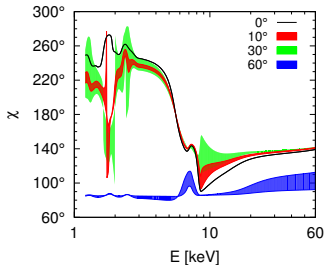
$a=0, \theta_0=60^\circ, h=15$



$a=1, \theta_0=60^\circ, h=3$



$a=1, \theta_0=60^\circ, h=15$



# Conclusions

- ▶ relativistic effects change considerably the polarization of the emitted radiation
- ▶ importance of the system geometry and physical properties of the system
- ▶ effects of “large-scale” corona (in XRBs) → Schnittman & Krolik (2010)
- ▶ the differences between results for various heights get smaller for polarisation degree
- ▶ the irregularities in the disc lower the polarisation degree and change the angle more with higher  $\delta_{\max}$
- ▶ the polarisation properties for  $\delta_{\max} \approx 60^\circ$  are similar for different heights and BH spins
- ▶ polarimetric observations could help to discriminate between competing scenarios and to determine the properties of the system  
→ inclination, BH spin, corona geometry