

Off-axis response of Compton and photoelectric polarimeters with a large field of view

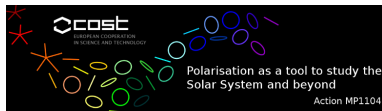
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**X-ray polarisation in astrophysics
-a window about to open?**

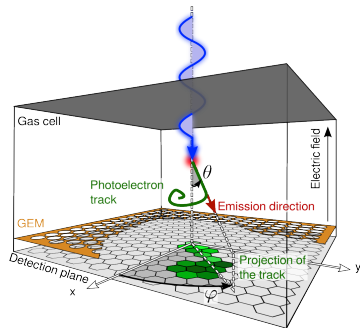
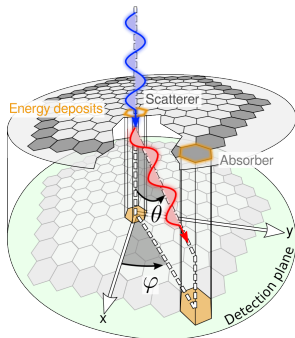
Stockholm, Sweden, 25-28 August 2014



- An analytical approach to calculate the modulation curve
- Difference between on-axis and off-axis photons
- A selection of the results
- Comparison with measurements
- Conclusions

How we build the modulation curve

- The modulation curve \mathcal{M} is just the number of counts produced in the direction φ



- Common instrumental designs are **not sensitive** to the angle ϑ
 - ➔ For example, Compton telescopes are an exception

How can this be expressed mathematically?

- The number of events produced in a certain direction is the **differential cross section** of the process:

$$\text{Photoelectric (K - shell)} : \frac{d\sigma_{\text{Ph}}}{d\Omega} = r_0^2 \alpha^4 Z^5 \left(\frac{m_e c^2}{E} \right)^{\frac{7}{2}} \frac{4\sqrt{2} \sin^2 \theta \cos^2 \phi}{(1 - \beta \cos \theta)^4}$$

$$\text{Compton (Klein - Nishina)} : \frac{d\sigma_{\text{Cm}}}{d\Omega} = \frac{1}{2} r_0^2 \frac{E'^2}{E^2} \left[\frac{E}{E'} + \frac{E'}{E} - 2 \sin^2 \theta \cos^2 \phi \right]$$

- Usually, the instruments are not sensitive to θ and then we have to **integrate** over this angle:

$$\mathcal{M}(\varphi) \propto \int_{\text{Min}}^{\text{Max}} \frac{d\sigma}{d\Omega} d\theta$$

How is this applied on-axis?

- We have to sum the number of events in the meridian slice φ

For completely polarized photons and a photoelectric polarimeter:

$$\mathcal{M}(\varphi) \propto \int_{\text{Min}}^{\text{Max}} \frac{\sin^2 \theta \cos^2 \phi}{(1 - \beta \cos \theta)^4} d\theta$$

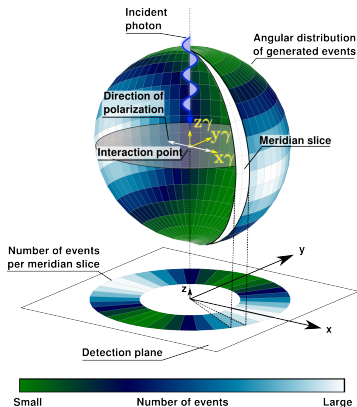
- On axis, this integral is trivial:

$$\mathcal{M}(\varphi) \propto \cos^2 \phi \underbrace{\left[\int_{\text{Min}}^{\text{Max}} \frac{\sin^2 \theta}{(1 - \beta \cos \theta)^4} d\theta \right]}_{\text{constant}=A}$$

In the following, $\text{Min}=0$ and $\text{Max}=\pi$.

- Taking into account the error in the measurements of the event direction:

$$\mathcal{M}(\varphi) = A \cos^2(\varphi - \varphi_0) + B$$

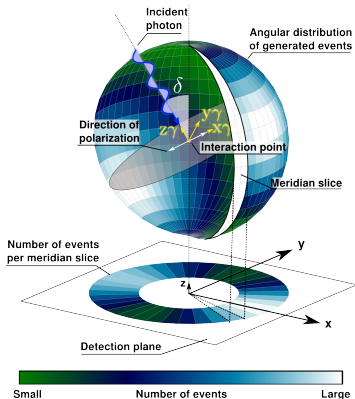
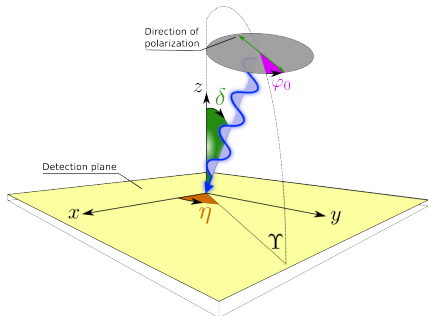


[Photoelectric, 100% polarized]

- The azimuthal and meridian directions are not decoupled as seen by in instrument
- The angular dependency of the events must be transformed in the instrument frame of reference

I used three rotations, representing

- δ : inclination
- η : azimuthal direction of the photons
- φ_0 : polarization angle



[Photoelectric, 100% polarized]

- There is no **conceptual** difference between photoelectric and Compton polarimeters
- Some algebra is required to arrive at the modulation curve
- In a simple configuration ($\eta = 0$ and $\varphi_0=0$), with a 1st order approximation in energy, for photoelectric polarimeters:

$$\begin{aligned}
 \mathcal{M}_{\text{Ph}} = f\mathcal{N}_{\text{tot}} & \left\{ \mathcal{P} \left[-\frac{9\beta \cos^2 \delta \sin \delta}{8} \cos^3 \varphi + \frac{\cos^2 \delta}{\pi} \cos^2 \varphi + \frac{3\beta (3\cos^2 \delta - 1) \sin \delta}{8} \cos \varphi + \frac{\sin^2 \delta}{2\pi} \right] + \right. \\
 & \left. + (1 - \mathcal{P}) \left[\frac{9\beta \sin^3 \delta}{16} \cos^3 \varphi - \frac{\sin^2 \delta}{2\pi} \cos^2 \varphi + \frac{3\beta(3\cos^2 \delta - 4) \sin \delta}{16} \cos \varphi + \frac{3 - \cos^2 \delta}{4\pi} \right] \right\} + \\
 & + \mathcal{N}_{\text{tot}} \frac{1 - f}{2\pi}
 \end{aligned}$$

\propto Pol. component
 \propto Unp. component
 \propto Instr. sensitivity

f : f-factor, conceptually similar to the modulation factor

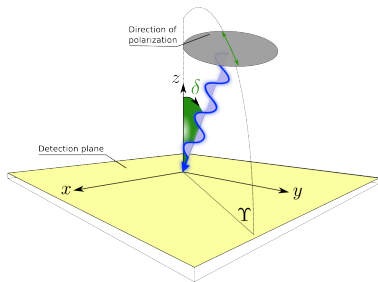
▶▶▶ In the following we will assume an **ideal instrument** with $f = 1$

\mathcal{P} : polarization degree

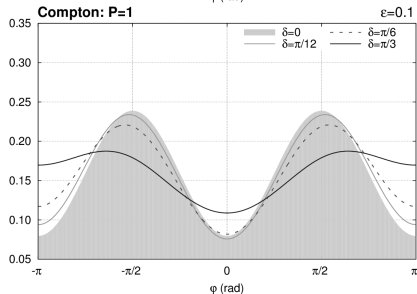
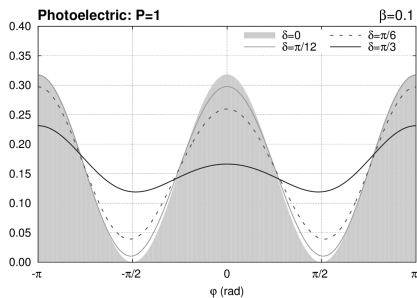
\mathcal{N}_{tot} : number of collected events

β : velocity of the photoelectron in c units, basically the energy of the photon

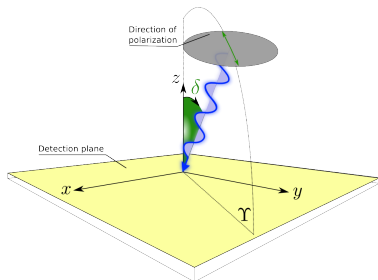
Increasing the inclination angle δ for completely polarized photons



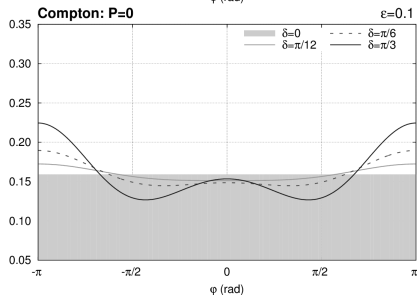
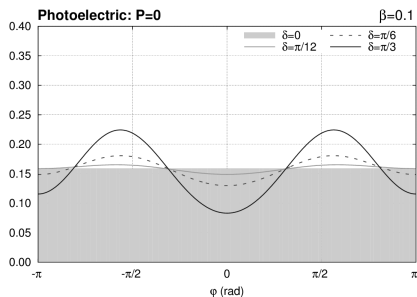
- Deviation from the cosine square behavior
 - Asymmetry of the two peaks
 - The modulation curve must be resolved over 360°
- Decreasing amplitude of the modulation with δ
- Increasing effects with the energy



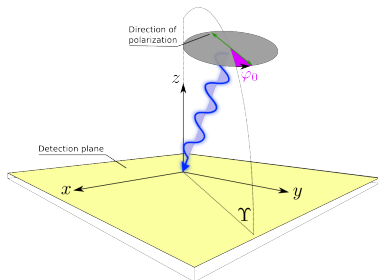
Increasing the inclination angle δ for completely unpolarized photons



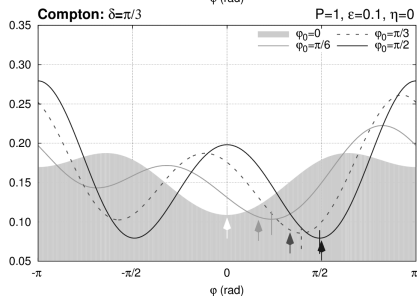
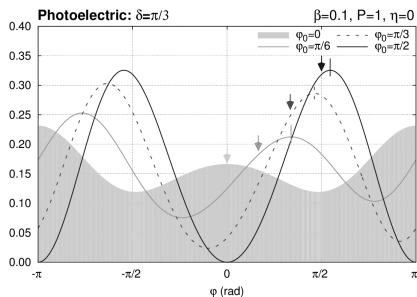
- A modulation is present also for off-axis unpolarized photons
 - In the low energy limit, this is exactly a **cosine square**
 - The amplitude increases with the inclination



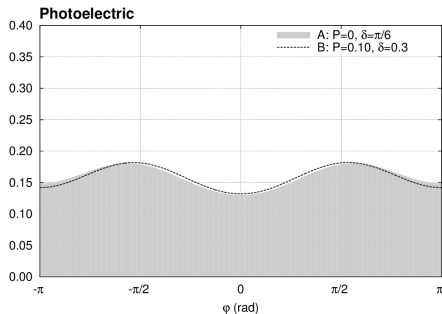
Changing the angle of polarization φ_0



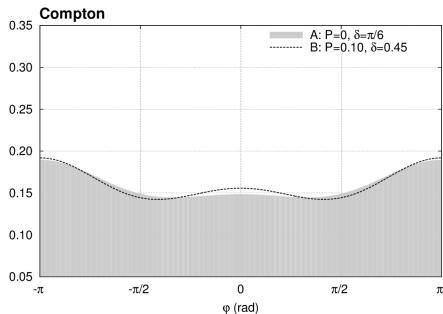
- The shape of the modulation curve depends on the angle of polarization
- The peak of the modulation does not correspond exactly to the polarization angle, even for 100% polarized photons



Comparison of 10% polarized and unpolarized photons with different inclination:



$$\Delta\delta = 13^\circ$$



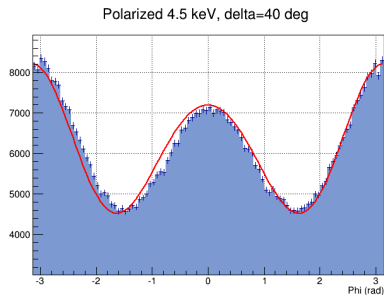
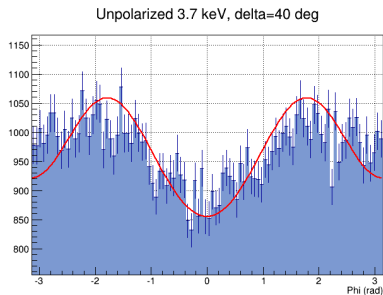
$$\Delta\delta = 4^\circ$$

A similar result holds true if we consider different **energies**.



The incident direction and energy of the photons **have to be known** to some degree

Photoelectric polarimeter (Gas Pixel Detector):



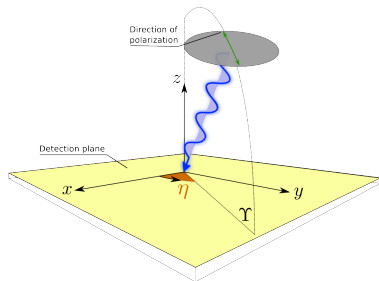
Red curves are not fit to data, but **theory predictions without any free parameter.**

- The modulation curve depends not only on the polarization, but also on the incident direction and energy in a **complex but predictable** way
 - ➡ The source position and the photon energy have to be known
- This is true for both usual “2-dimensional” photoelectric and Compton polarimeters
 - ➡ The only way of avoiding this is 3-dimensional polarimeters, like Compton telescopes, with a **complete** reconstruction of the event direction.

[see Muleri 2014, The Astrophysical Journal, 782:28 for more details]

Backup slides

Changing the azimuthal direction of the photons η



- The modulation curve just shift of η
 - This is true for the modulation produced by both polarized and unpolarized photons

